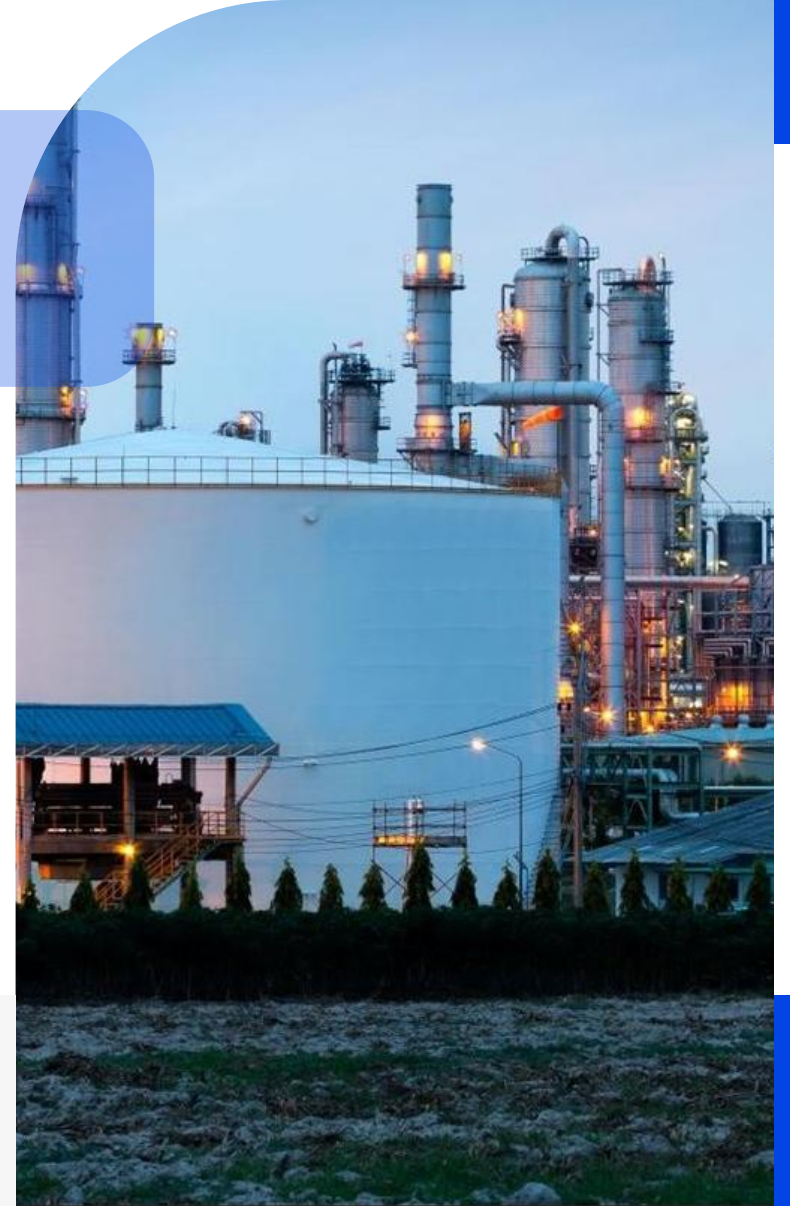


Thermodynamic Properties of Fluids

Termodinamika Teknik Kimia I



Pendahuluan

- **Kaidah Fase Gibbs:** jika sejumlah sifat intensif telah ditetapkan, maka sifat-sifat intensif yang lain akan tertentu pula.
- Nilai numerik suatu sifat termodinamika penting untuk menghitung a.l. : **kerja dan panas dalam suatu proses, kesetimbangan fase/ reaksi kimia.**
- *Perlu adanya persamaan hubungan antar sifat intensif, juga hubungan sifat intensif dengan persamaan keadaan (EOS) atau besaran-besaran yang mudah terukur (P , V atau T)*

**Besaran terukur
(P, V, T)**

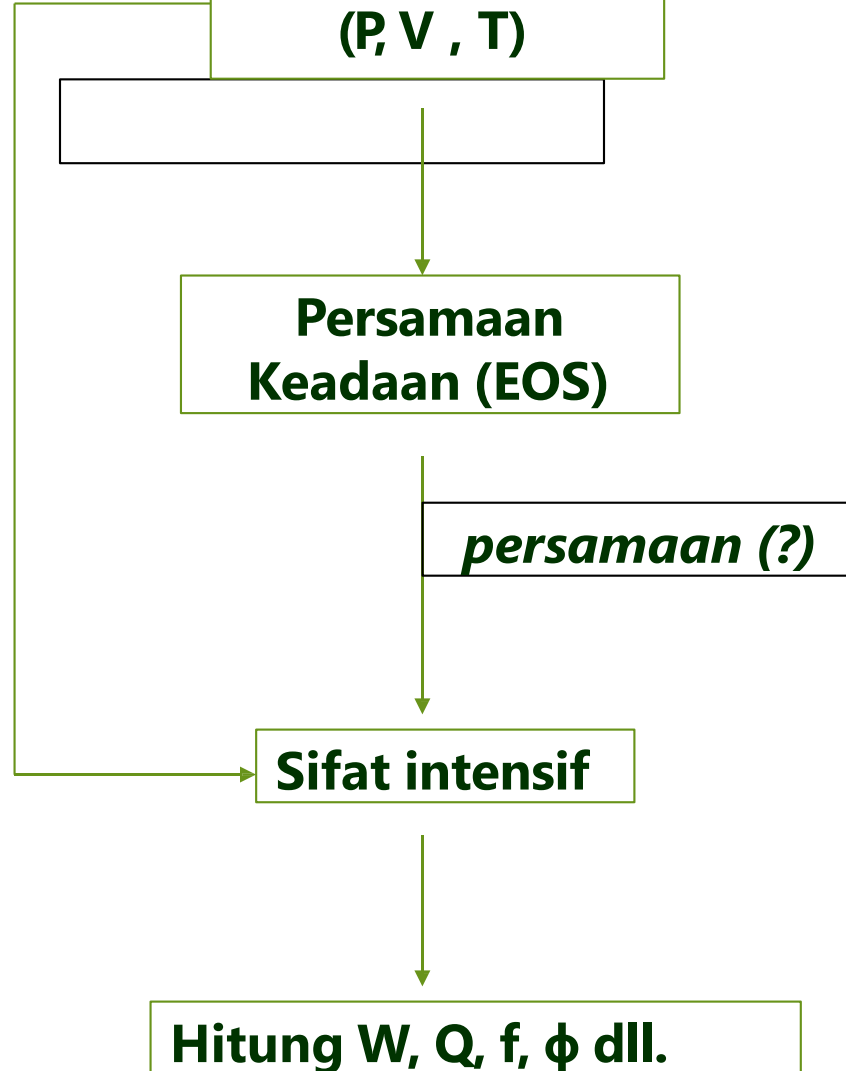
persamaan (?)

**Persamaan
Keadaan (EOS)**

persamaan (?)

Sifat intensif

Hitung W, Q, f, ϕ dll.



Hubungan antar Sifat untuk Fase-fase Homogen

Hukum I untuk sistem tertutup:

$$d(nU) = dQ + dW$$

khusus proses reversibel

$$d(nU) = dQ_{rev} + dW_{rev}$$

$$dW_{rev} = -Pd(nV)$$

$$dQ_{rev} = Td(nS)$$

Hukum II

$$d(nU) = Td(nS) - Pd(nV)$$

- Persamaan tersebut hanya mengandung sifat-sifat sistem.
- Dapat diterapkan pada sistem tertutup dengan massa konstan, baik pada proses reversible atau irreversible.
- Perubahan sifat (properties change) di antara keadaan-keadaan yang setimbang.
- Sistem bisa berupa fase homogen atau heterogen, inert atau terlibat reaksi kimia.

Persamaan *) mengandung sifat-sifat termodinamika yg utama, yaitu:

$$d(nU) = Td(nS) - Pd(nV) \quad *)$$

$$P, V, T, U, \text{ and } S$$

Entalpi:

$$H \equiv U + PV$$

Energi (bebas) Helmholtz:

$$A \equiv U - TS$$

Energi (bebas) Gibbs:

$$G \equiv H - TS$$

Untuk satu mol fluida homogen dengan komposisi konstan:

$$dU = TdS - PdV$$

$$dH = TdS + VdP$$

$$dA = -PdV - SdT$$

$$dG = VdP - SdT$$

PD
eksak

$$\left(\frac{\partial T}{\partial V} \right)_S = - \left(\frac{\partial P}{\partial S} \right)_V$$

$$\left(\frac{\partial T}{\partial P} \right)_S = \left(\frac{\partial V}{\partial S} \right)_P$$

$$\left(\frac{\partial P}{\partial T} \right)_V = \left(\frac{\partial S}{\partial V} \right)_T$$

$$\left(\frac{\partial V}{\partial T} \right)_P = - \left(\frac{\partial S}{\partial P} \right)_T$$

Maxwell's
equations

Perubahan Entalpi $f(P,T)$

$$dH = \left(\frac{\partial H}{\partial T} \right)_P dT + \left(\frac{\partial H}{\partial P} \right)_T dP$$

$$dH = TdS + VdP$$

$$\left(\frac{\partial H}{\partial T} \right)_P = C_P$$

$$\left(\frac{\partial H}{\partial P} \right)_T = T \left(\frac{\partial S}{\partial P} \right)_T + V$$

$$\left(\frac{\partial H}{\partial T} \right)_P = T \left(\frac{\partial S}{\partial T} \right)_P$$

$$dH = C_P dT + \left(T \left(\frac{\partial S}{\partial P} \right)_T + V \right) dP$$

$$\left(\frac{\partial V}{\partial T} \right)_P = - \left(\frac{\partial S}{\partial P} \right)_T$$

$$dH = C_P dT + \left(V - T \left(\frac{\partial V}{\partial T} \right)_P \right) dP$$

Perubahan Entropi $f(P,T)$

$$dS = \left(\frac{\partial S}{\partial T} \right)_P dT + \left(\frac{\partial S}{\partial P} \right)_T dP$$

$$\left(\frac{\partial S}{\partial T} \right)_P = \frac{C_P}{T}$$

$$\left(\frac{\partial V}{\partial T} \right)_P = - \left(\frac{\partial S}{\partial P} \right)_T$$

$$dS = C_P \frac{dT}{T} - \left(\frac{\partial V}{\partial T} \right)_P dP$$

Perubahan Energi Dalam $f(P,T)$

$$H \equiv U + PV$$

Pers.
6.19

$$\left(\frac{\partial H}{\partial P}\right)_T = \left(\frac{\partial U}{\partial P}\right)_T + P\left(\frac{\partial V}{\partial P}\right)_T + V$$

$$\left(\frac{\partial U}{\partial P}\right)_T = -T\left(\frac{\partial V}{\partial T}\right)_P - P\left(\frac{\partial V}{\partial P}\right)_T$$

Perubahan Entalpi & Entropi Gas Ideal

$$PV^{ig} = RT$$

$$\left(\frac{\partial V^{ig}}{\partial T} \right)_P = \frac{R}{P}$$

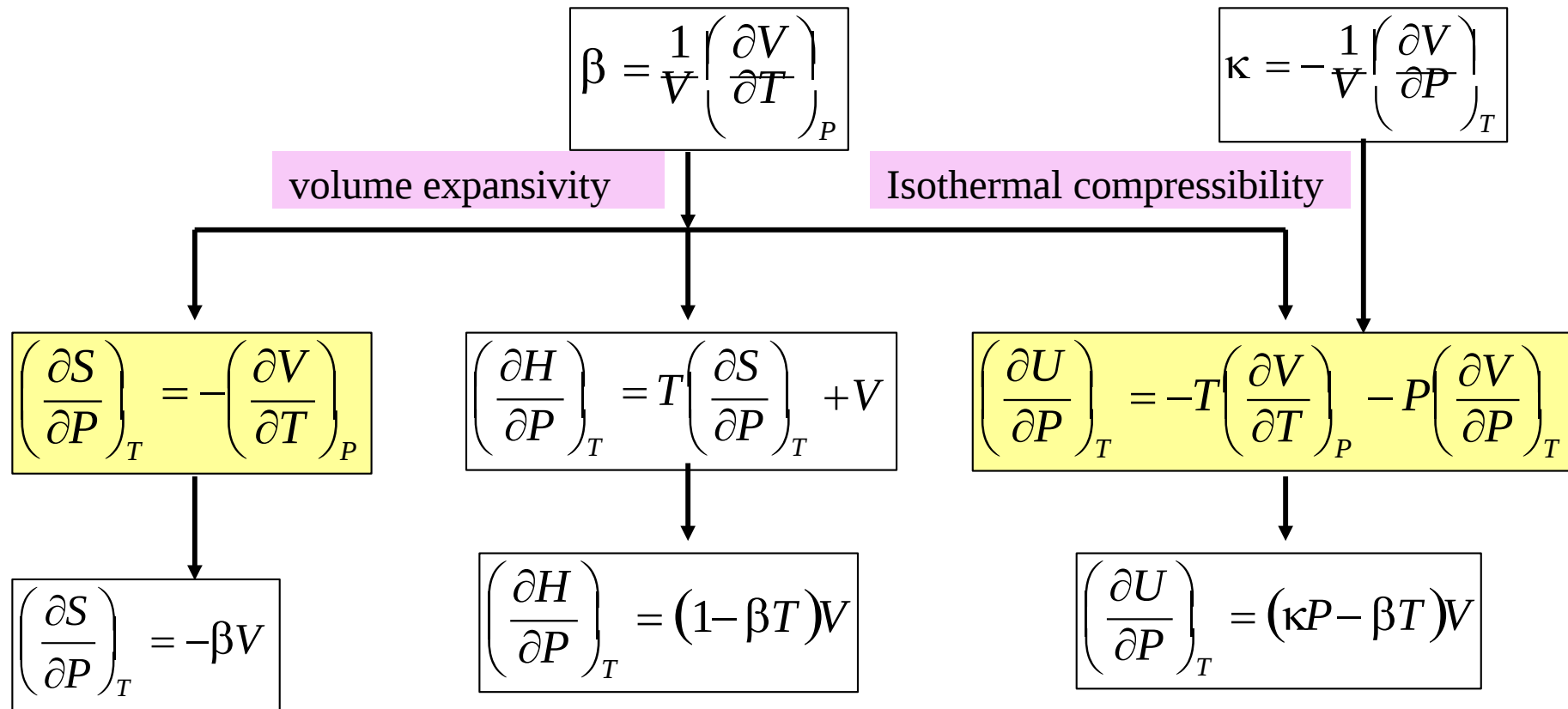
$$dH = C_P dT + \left(V - T \left(\frac{\partial V}{\partial T} \right)_P \right) dP$$

$$dH^{ig} = C_P^{ig} dT$$

$$dS = C_P \frac{dT}{T} - \left(\frac{\partial V}{\partial T} \right)_P dP$$

$$dS^{ig} = C_P^{ig} \frac{dT}{T} - R \frac{dP}{P}$$

Bentuk Alternatif untuk Cairan



Bentuk Alternatif untuk Cairan (lanjutan)

$$\left(\frac{\partial V}{\partial T}\right)_P = \beta V$$

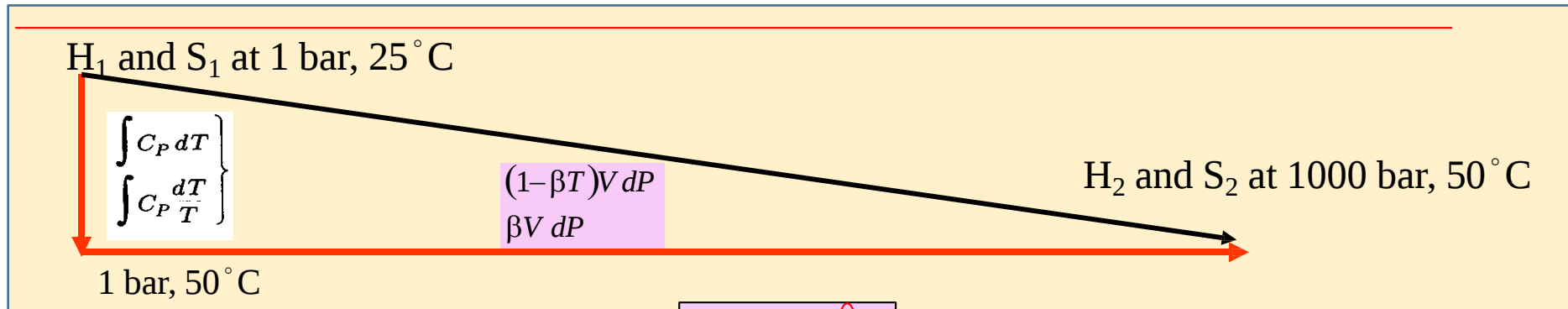
$$dH = C_P dT + \left[V - T \left(\frac{\partial V}{\partial T}\right)_P \right] dP$$

$$dH = C_P dT + (1 - \beta T) V dP$$

$$dS = C_P \frac{dT}{T} - \left(\frac{\partial V}{\partial T}\right)_P dP$$

$$dS = C_P \frac{dT}{T} - \beta V dP$$

Determine the enthalpy and entropy changes of liquid water for a change of state from 1 bar and 25 °C to 1000 bar and 50 °C. The data for water are given.



$$dH = C_P dT + \left(V - T \left(\frac{\partial V}{\partial T} \right)_P \right) dP$$

$$\left(\frac{\partial V}{\partial T} \right)_P = \beta V$$

volume expansivity

$$dH = C_P dT + (1 - \beta T) V dP$$

$$\Delta H = \langle C_P \rangle (T_2 - T_1) + (1 - \langle \beta \rangle T_2) \langle V \rangle (P_2 - P_1)$$

$$\Delta H = 75.310(323.15 - 298.15) + \frac{(1 - (513 \times 10^{-6})(323.15))(18.204)(1000 - 1)}{10} = 3400 \text{ J / mol}$$

$$dS = C_P \frac{dT}{T} - \left(\frac{\partial V}{\partial T} \right)_P dP$$

$$\Delta S = \langle C_P \rangle \ln \frac{T_2}{T_1} - \langle \beta \rangle \langle V \rangle (P_2 - P_1)$$

$$\Delta S = 75.310 \ln \frac{323.15}{298.15} - \frac{(513 \times 10^{-6})(18.204)(1000 - 1)}{10} = 5.13 \text{ J / mol}$$

Energi Dalam & Entropi f (V,T)

$$dU = \left(\frac{\partial U}{\partial T} \right)_V dT + \left(\frac{\partial U}{\partial V} \right)_T dV$$

$$\left(\frac{\partial U}{\partial T} \right)_V = C_V$$

$$\left(\frac{\partial U}{\partial V} \right)_T = T \left(\frac{\partial S}{\partial V} \right)_T - P$$

$$dU = C_V dT + \left[T \left(\frac{\partial S}{\partial V} \right)_T - P \right] dV$$

$$\left(\frac{\partial S}{\partial V} \right)_T = - \left(\frac{\partial P}{\partial T} \right)_V$$

$$dU = C_V dT + \left[T \left(\frac{\partial P}{\partial T} \right)_V - P \right] dV$$

$$dS = \left(\frac{\partial S}{\partial T} \right)_V dT + \left(\frac{\partial S}{\partial V} \right)_T dV$$

$$\left(\frac{\partial U}{\partial T} \right)_V = T \left(\frac{\partial S}{\partial T} \right)_V$$

$$\left(\frac{\partial S}{\partial V} \right)_T = - \left(\frac{\partial P}{\partial T} \right)_V$$

$$\left(\frac{\partial S}{\partial T} \right)_V = \frac{C_V}{T}$$

$$dS = C_V \frac{dT}{T} + \left(\frac{\partial P}{\partial T} \right)_V dV$$

Bentuk alternatif

$$\left(\frac{\partial P}{\partial T}\right)_V = \frac{\beta}{\kappa}$$

$$dU = C_V dT + \left(T \left(\frac{\partial P}{\partial T}\right)_V - P \right) dV$$

$$dU = C_V dT + \left(\frac{\beta}{\kappa} T - P \right) dV$$

$$dS = C_V \frac{dT}{T} + \left(\frac{\partial P}{\partial T}\right)_V dV$$

$$dS = \frac{C_V}{T} dT + \frac{\beta}{\kappa} dV$$

Gibbs energy $G = G(P, T)$

$$dG = VdP - SdT$$

Thermodynamic property of great potential utility

$$d\left(\frac{G}{RT}\right) \equiv \frac{1}{-} dG - \frac{G}{RT^2} dT$$

$$G \equiv H - TS$$

$$dG = VdP - SdT$$

$$d\left(\frac{G}{RT}\right) \equiv \frac{V}{RT} dP - \frac{H}{RT^2} dT$$

$$d\left(\frac{G}{RT}\right) = \left(\frac{\partial(G/RT)}{\partial P}\right)_T dP + \left(\frac{\partial(G/RT)}{\partial T}\right)_P dT$$

$$\frac{V}{RT} = \left(\frac{\partial(G/RT)}{\partial P}\right)_T$$

$$\frac{H}{RT} = -T \left(\frac{\partial(G/RT)}{\partial T}\right)_P$$

$$G/RT = g(P, T)$$

The Gibbs energy serves as a generating function for the other thermodynamic properties.

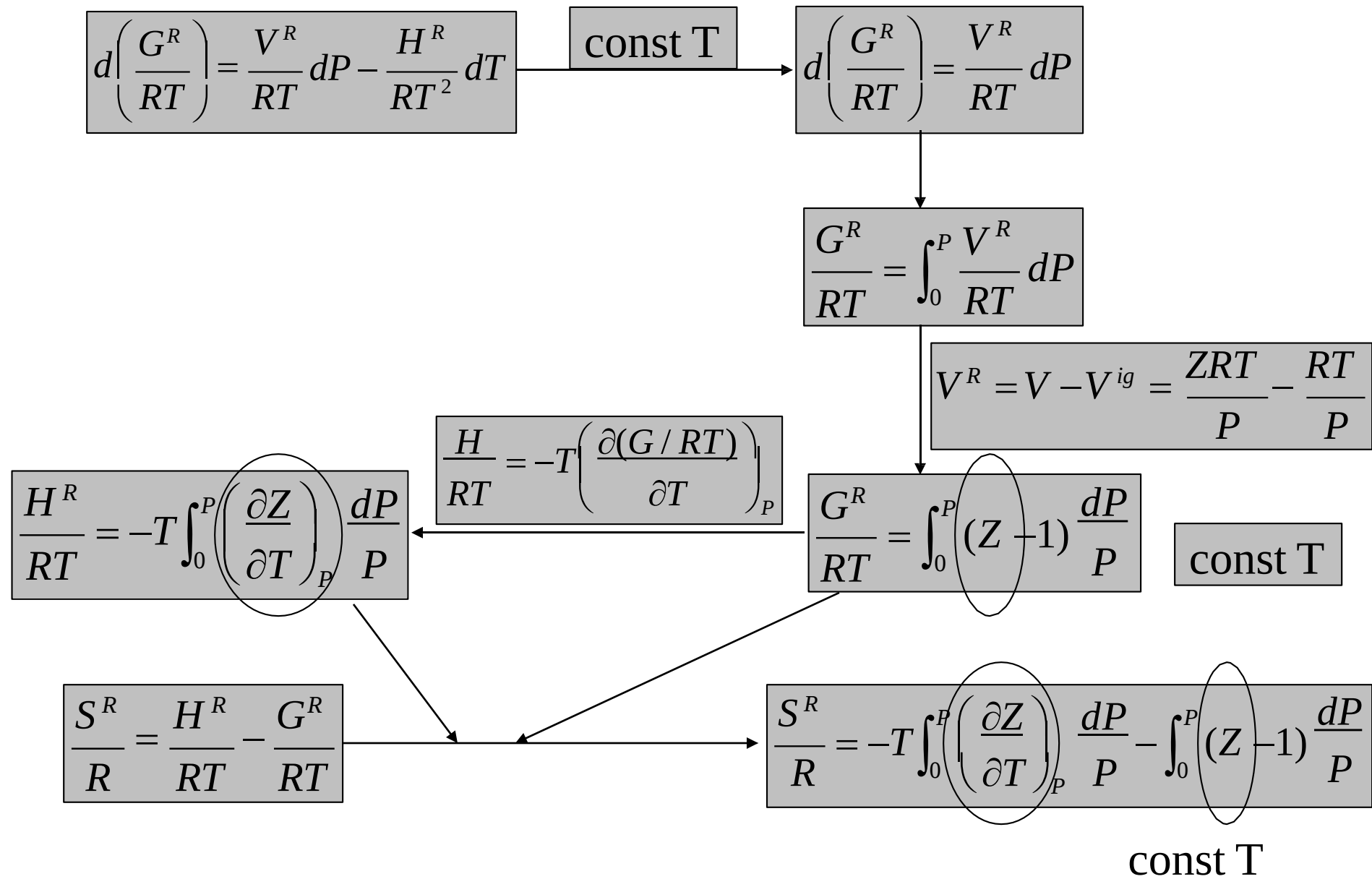
Residual properties

- The definition for the generic residual property:

$$M^R \equiv M - M^{ig}$$

- M and M^{ig} are the actual and ideal-gas properties,
 - respectively.
 - M is the molar value of any extensive thermodynamic properties, e.g., V , U , H , S , or G .
- The residual Gibbs energy serves as a generating function for the other residual properties:

$$d\left(\frac{G^R}{RT}\right) = \frac{V^R}{RT} dP - \frac{H^R}{RT^2} dT$$



$Z = PV/RT$: experimental measurement .Given PVT data or an appropriate equation of state, we can evaluate H^R and S^R and hence all other residual properties.

$$\begin{aligned}
H &= H^{ig} + H^R = \left\{ H_0^{ig} + \int_{T_0}^T C_P^{ig} dT \right\} + \left\{ -RT^2 \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} \right\} \\
&= \left\{ H_0^{ig} + R \times ICPH(T_0, T; A, B, C, D) \right\} + \left\{ -RT^2 \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} \right\} \\
&= \left\{ H_0^{ig} + \langle C_P^{ig} \rangle_H (T - T_0) \right\} + \left\{ -RT^2 \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} \right\} \\
&= \left\{ H_0^{ig} + R \times MCPH(T_0, T; A, B, C, D)(T - T_0) \right\} + \left\{ -RT^2 \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} \right\}
\end{aligned}$$

$$\begin{aligned}
S &= S^{ig} + S^R = \left\{ S_0^{ig} + \int_{T_0}^T C_P^{ig} \frac{dT}{T} - R \ln \frac{P}{P_0} \right\} + R \left\{ -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} - \int_0^P (Z - 1) \frac{dP}{P} \right\} \\
&= \left\{ S_0^{ig} + R \times ICPS(T_0, T; A, B, C, D) - R \ln \frac{P}{P_0} \right\} + R \left\{ -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} - \int_0^P (Z - 1) \frac{dP}{P} \right\} \\
&= \left\{ S_0^{ig} + \langle C_P^{ig} \rangle_S \ln \frac{T}{T_0} - R \ln \frac{P}{P_0} \right\} + R \left\{ -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} - \int_0^P (Z - 1) \frac{dP}{P} \right\} \\
&= \left\{ S_0^{ig} + R \times MCPS(T_0, T; A, B, C, D) \ln \frac{T}{T_0} - R \ln \frac{P}{P_0} \right\} + R \left\{ -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} - \int_0^P (Z - 1) \frac{dP}{P} \right\}
\end{aligned}$$

Calculate the enthalpy and entropy of saturated isobutane vapor at 360 K from the following information: (1) compressibility-factor for isobutane vapor; (2) the vapor pressure of isobutane at 360 K is 15.41 bar; (3) at 300K and 1 bar, (4) the ideal-gas heat capacity of isobutane vapor:

$$H_0^{ig} = 18115 \text{ J / mol } S_0^{ig} = 295.976 \text{ J / mol} \cdot \text{K}$$

$$C_P^{ig} / R = 1.7765 + 33.037 \times 10^{-3} T$$

$$\boxed{\frac{H^R}{RT} = -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P}} \quad \boxed{\frac{S^R}{R} = -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} - \int_0^P (Z - 1) \frac{dP}{P}}$$

Graphical integration requires plots of $\left(\frac{\partial Z}{\partial T} \right)_P / P$ and $(Z-1)/P$ vs. P .

$$\boxed{\frac{H^R}{RT} = -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} = -(360)(26.37 \times 10^{-4}) = -0.9493} \longrightarrow \boxed{H^R = -2841.3 \text{ J / mol}}$$

$$\boxed{\frac{S^R}{R} = -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} - \int_0^P (Z - 1) \frac{dP}{P} = -0.9493 - (-0.02596) = -0.6897} \longrightarrow \boxed{S^R = -5.734 \text{ J / mol} \cdot \text{K}}$$

$$\boxed{H = \{H_0^{ig} + R \times ICPH(300, 360; 1.7765, 33.037E-3, 0.0, 0.0)\} + H^R = 21598.5 \frac{\text{J}}{\text{mol}}}$$

$$\boxed{S = \{S_0^{ig} + R \times ICPS(300, 360; 1.7765, 33.037E-3, 0.0, 0.0)\} - R \ln \frac{15.41}{1} + H^R = 286.676 \frac{\text{J}}{\text{mol} \cdot \text{K}}}$$

Residual properties by equations of state

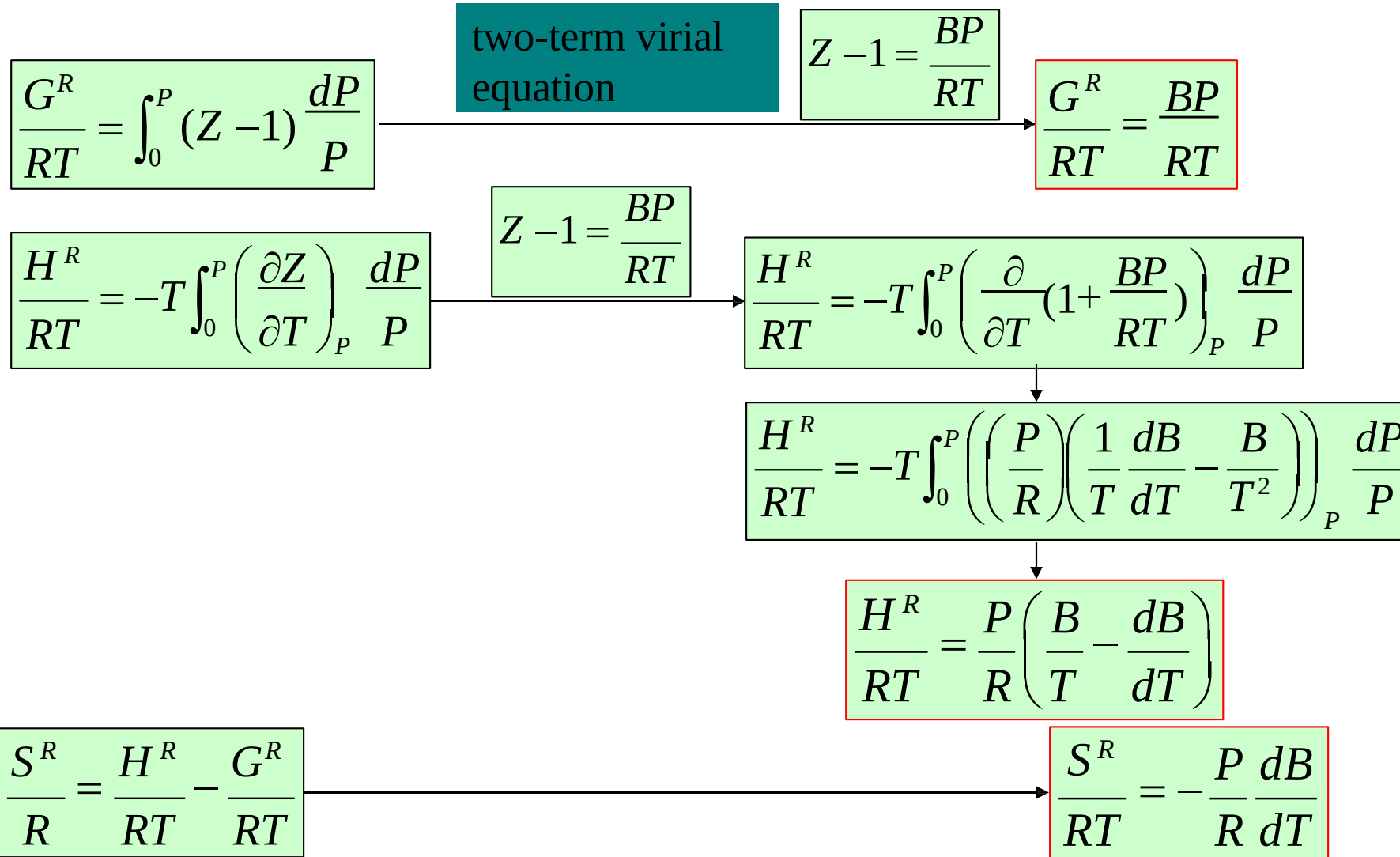
- If $Z = f(P, T)$:

$$\frac{H^R}{RT} = -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P}$$

$$\frac{S^R}{R} = -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} - \int_0^P (Z - 1) \frac{dP}{P}$$

- The calculation of residual properties for gases and vapors through use of the virial equations and cubic equation of state.

Use of the virial equation of state



Three-term virial equation

$$Z - 1 = B\rho + C\rho^2$$

$$P = Z\rho RT$$

$$\frac{G^R}{RT} = \int_0^P (Z - 1) \frac{dP}{P}$$

$$\frac{G^R}{RT} = 2B\rho + \frac{3}{2}C\rho^2 - \ln Z$$

$$\frac{H^R}{RT} = -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P}$$

$$\frac{H^R}{RT} = T \left[\left(\frac{B}{T} - \frac{dB}{dT} \right) \rho + \left(\frac{C}{T} - \frac{1}{2} \frac{dC}{dT} \right) \rho^2 \right]$$

$$\frac{S^R}{R} = \frac{H^R}{RT} - \frac{G^R}{RT}$$

Application up to moderate pressure

USE OF THE CUBIC EQUATION OF STATE

- The generic cubic equation of state:

$$P = \frac{RT}{V-b} - \frac{a(T)}{(V+\epsilon b)(V+\sigma b)}$$

$$q \equiv \frac{a(T)}{bRT}$$

$$Z = \frac{1}{1-\rho b} - q \frac{\rho b}{(1+\epsilon \rho b)(1+\sigma \rho b)}$$

$$\frac{H^R}{RT} = Z - 1 + \left[\frac{d \ln \alpha(T_r)}{d \ln T_r} - 1 \right] qI$$

$$\frac{S^R}{R} = \ln(Z - \beta) + \left[\frac{d \ln \alpha(T_r)}{d \ln T_r} \right] qI$$

$$\beta \equiv \frac{bP}{RT}$$

$$\int_0^\rho (Z-1) \frac{d\rho}{\rho} = -\ln(1-\rho b) - qI$$

$$\int_0^\rho \left(\frac{\partial Z}{\partial T} \right)_\rho \frac{d\rho}{\rho} = - \frac{dq}{dT} I$$

$$I \equiv \int_0^\rho \frac{d(\rho b)}{(1+\epsilon \rho b)(1+\sigma \rho b)} \quad \text{const } T$$

Find values for the residual enthalpy H^R and the residual entropy S^R for n-butane gas at 500 K and 50 bar as given by Redlich/Kwong equation.

$$T_r = \frac{500}{425.1} = 1.176 \quad P_r = \frac{50}{37.96} = 1.317 \quad \beta = \Omega \frac{P_r}{T_r} = 0.09703 \quad q = \frac{\Psi \alpha(T_r)}{\Omega T_r} = 3.8689$$

$$Z = 1 + \beta - q\beta \frac{Z - \beta}{(Z + \varepsilon\beta)(Z + \sigma\beta)} \quad \begin{matrix} \varepsilon = 0 \\ \sigma = 1 \end{matrix} \rightarrow Z = 0.685$$

$$I \equiv \int_0^{\rho} \frac{d(\rho b)}{(1 + \varepsilon \rho b)(1 + \sigma \rho b)} \quad \text{const } T \rightarrow I = \frac{1}{\sigma - \varepsilon} \ln \left(\frac{Z + \sigma\beta}{Z + \varepsilon\beta} \right) = 0.13247$$

$$\frac{H^R}{RT} = Z - 1 + \left[\frac{d \ln \alpha(T_r)}{d \ln T_r} - 1 \right] qI \rightarrow H^R = -4505 \frac{J}{mol}$$

$$\frac{S^R}{R} = \ln(Z - \beta) + \left[\frac{d \ln \alpha(T_r)}{d \ln T_r} \right] qI \rightarrow S^R = -6.546 \frac{J}{mol \cdot K}$$

Two-phase systems

- Whenever a phase transition at constant temperature and pressure occurs,
 - The molar or specific volume, internal energy, enthalpy, and entropy changes abruptly. The exception is the molar or specific Gibbs energy, which for a pure species does not change during a phase transition.
 - For two phases α and β of a pure species coexisting at equilibrium:

$$G^{\alpha} = G^{\beta}$$

where G^{α} and G^{β} are the molar or specific Gibbs energies of the individual phases

$$G^{\alpha} = G^{\beta} \longrightarrow dG^{\alpha} = dG^{\beta}$$

$$dG = VdP - SdT$$

$$V^{\alpha} dP^{sat} - S^{\alpha} dT = V^{\beta} dP^{sat} - S^{\beta} dT$$

$$\frac{dP^{sat}}{dT} = \frac{S^{\alpha} - S^{\beta}}{V^{\alpha} - V^{\beta}} = \frac{\Delta S^{\alpha\beta}}{\Delta V^{\alpha\beta}}$$

$$dH = TdS + VdP$$

$$\Delta H^{\alpha\beta} = T\Delta S^{\alpha\beta}$$

The latent heat of phase transition

The Clapeyron equation

$$\frac{dP^{sat}}{dT} = \frac{\Delta H^{\alpha\beta}}{T\Delta V^{\alpha\beta}}$$

$$\Delta V^{\alpha\beta} = V^v = \frac{RT}{P^{sat}}$$

Ideal gas, and $V^l \ll V^g$

$$\Delta H^{lv} = -R \frac{d \ln P^{sat}}{d(1/T)}$$

The Clausius/Clapeyron equation

$$\frac{dP^{sat}}{dT} = \frac{\Delta H^{\alpha\beta}}{T\Delta V^{\alpha\beta}}$$

The Clapeyron equation: a vital connection between the properties of different phases.

$$\longrightarrow P = f(T)$$

$$\Delta H^{lv} = -R \frac{d \ln P^{sat}}{d(1/T)}$$

The Clausius/Clapeyron equation; $\ln P$ vs $1/T \rightarrow$ linier

$$\ln P^{sat} = A - \frac{B}{T}$$

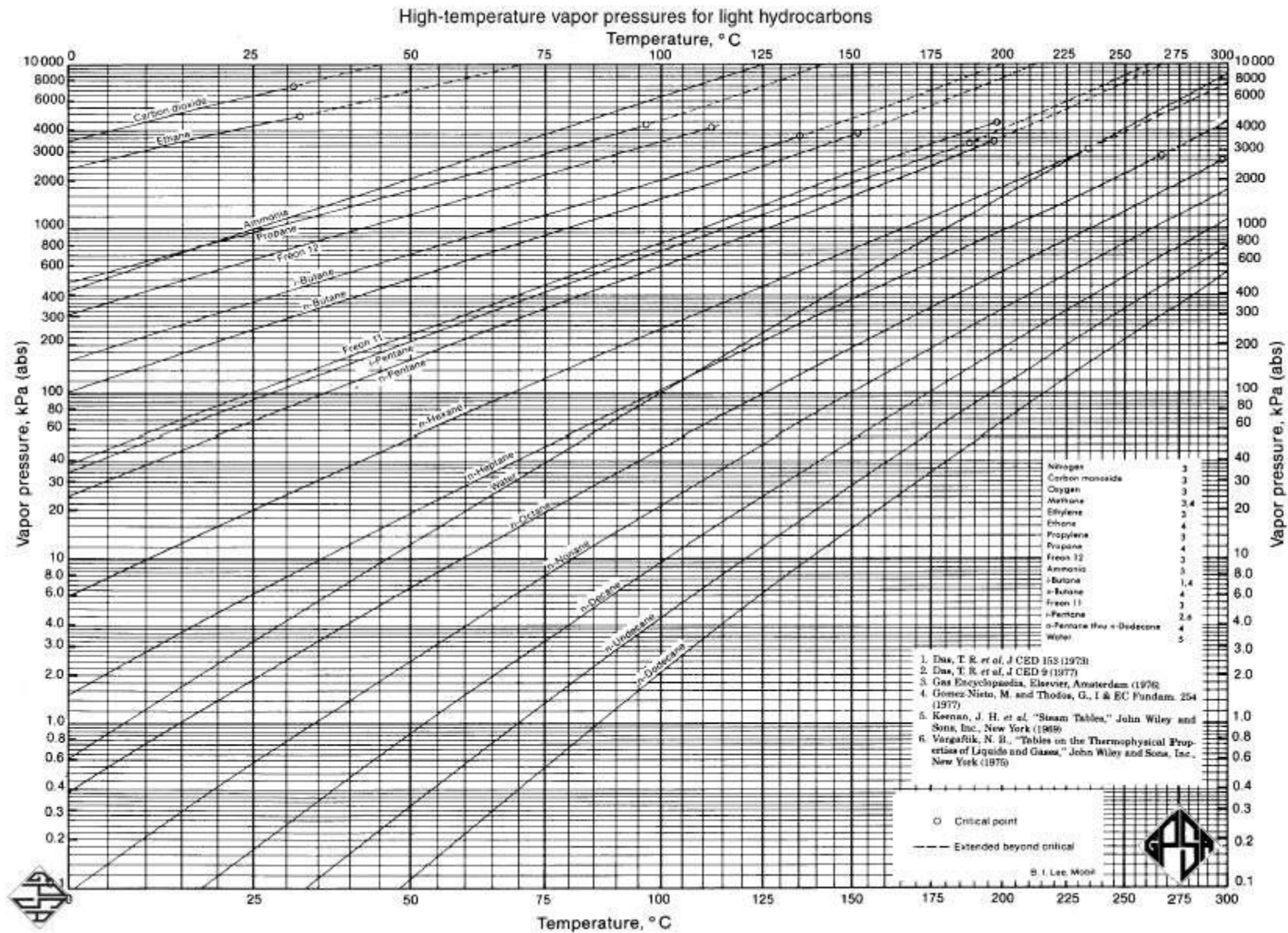
For the entire temperature range from the triple point to the critical point

$$\ln P^{sat} = A - \frac{B}{T + C}$$

The Antoine equation, for a specific temperature range

$$\left\{ \begin{array}{l} \ln P_r^{sat} = \frac{A\tau + B\tau^{1.5} + C\tau^3 + D\tau^6}{1-\tau} \\ \tau = 1 - T_r \end{array} \right.$$

The Wagner equation, over a wide temperature range.



High-Temperature Vapor Pressures for Light Hydrocarbons

FIG. 23-20

Two-phase liquid/ vapor systems

- When a system consists of saturated-liquid and saturated-vapor phases coexisting in equilibrium, the total value of any extensive property of the two-phase system is the sum of the total properties of the phases:

- M represents V, U, H, S, etc.

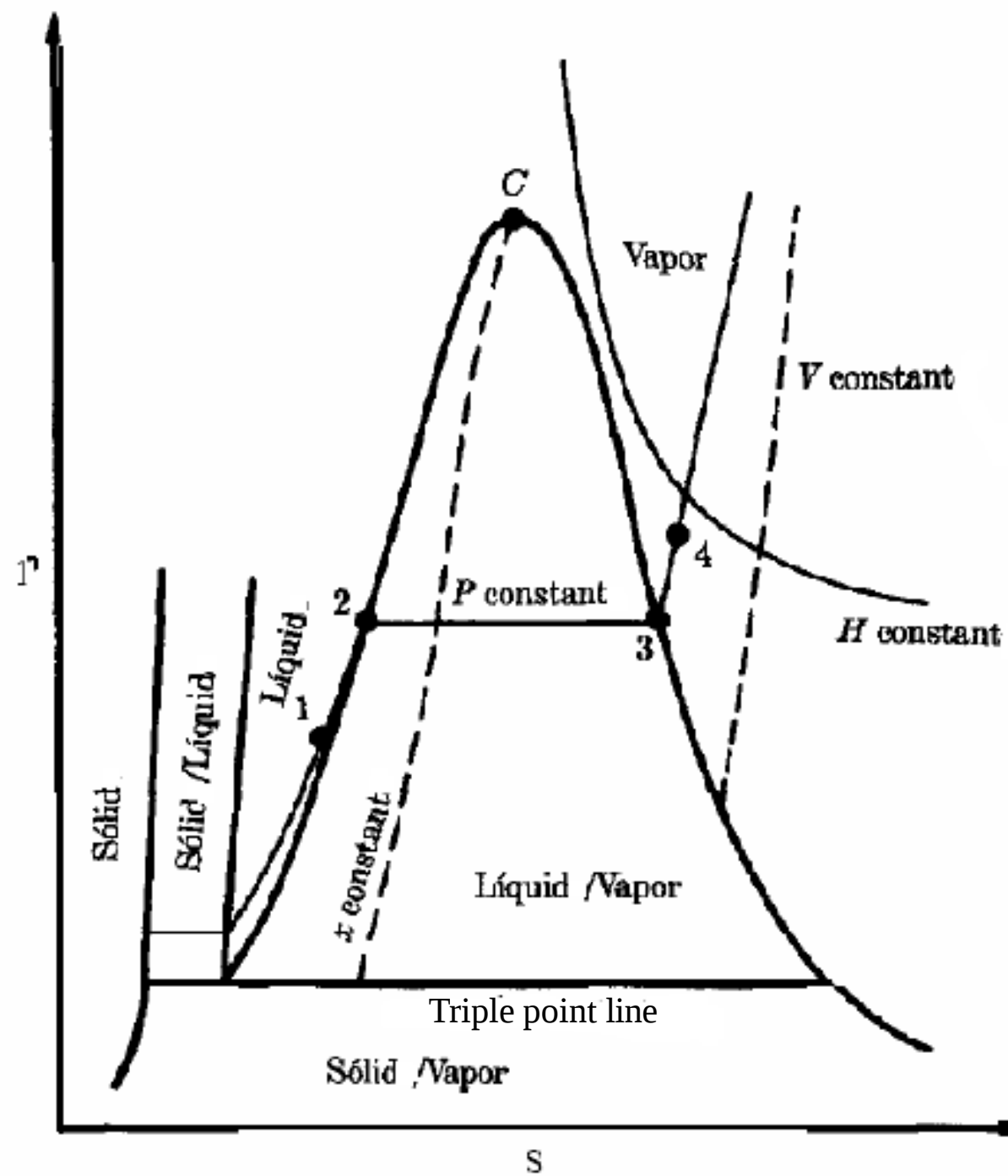
$$M = (1 - x^v)M^l + x^v M^v$$

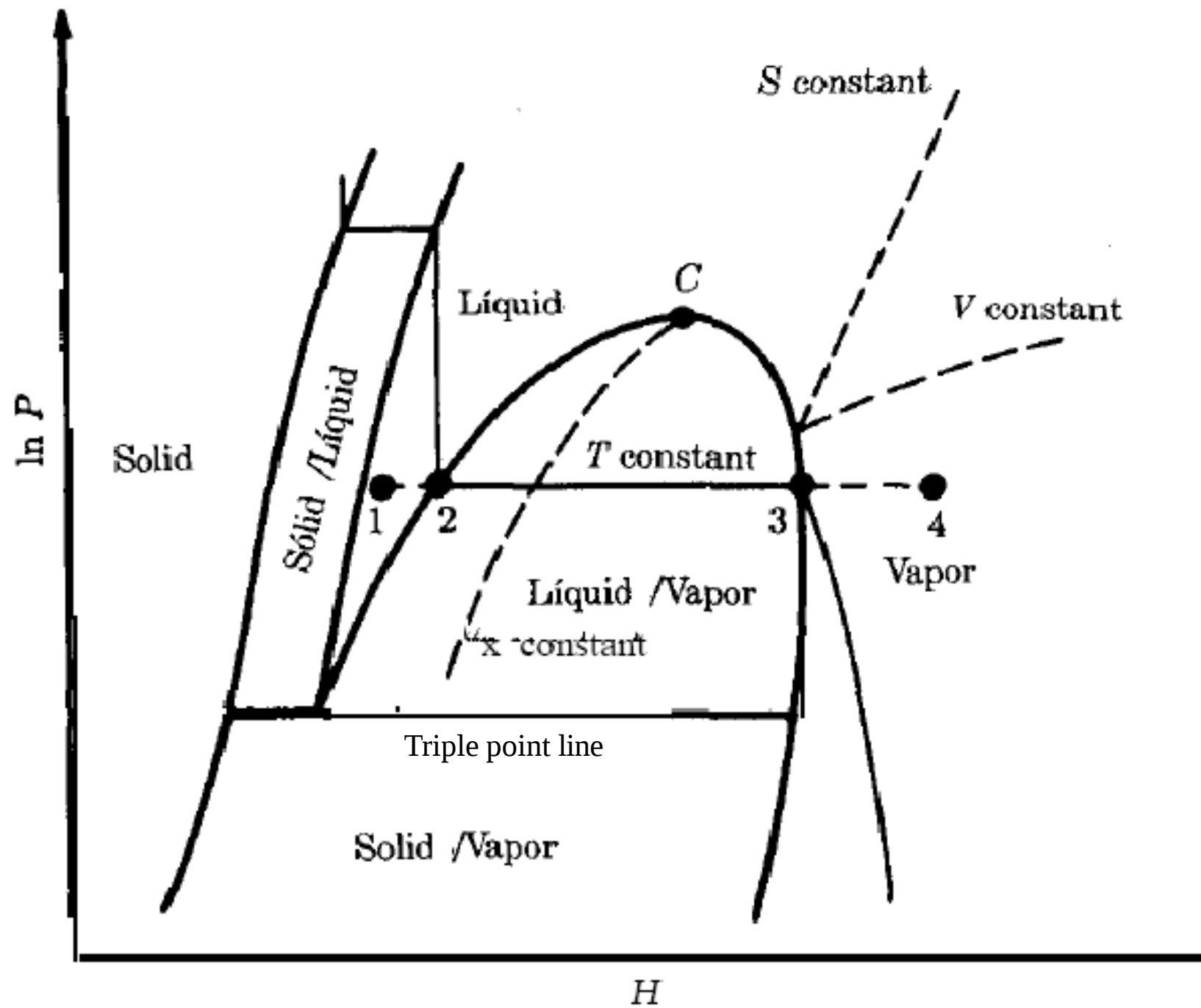
- e.g.,

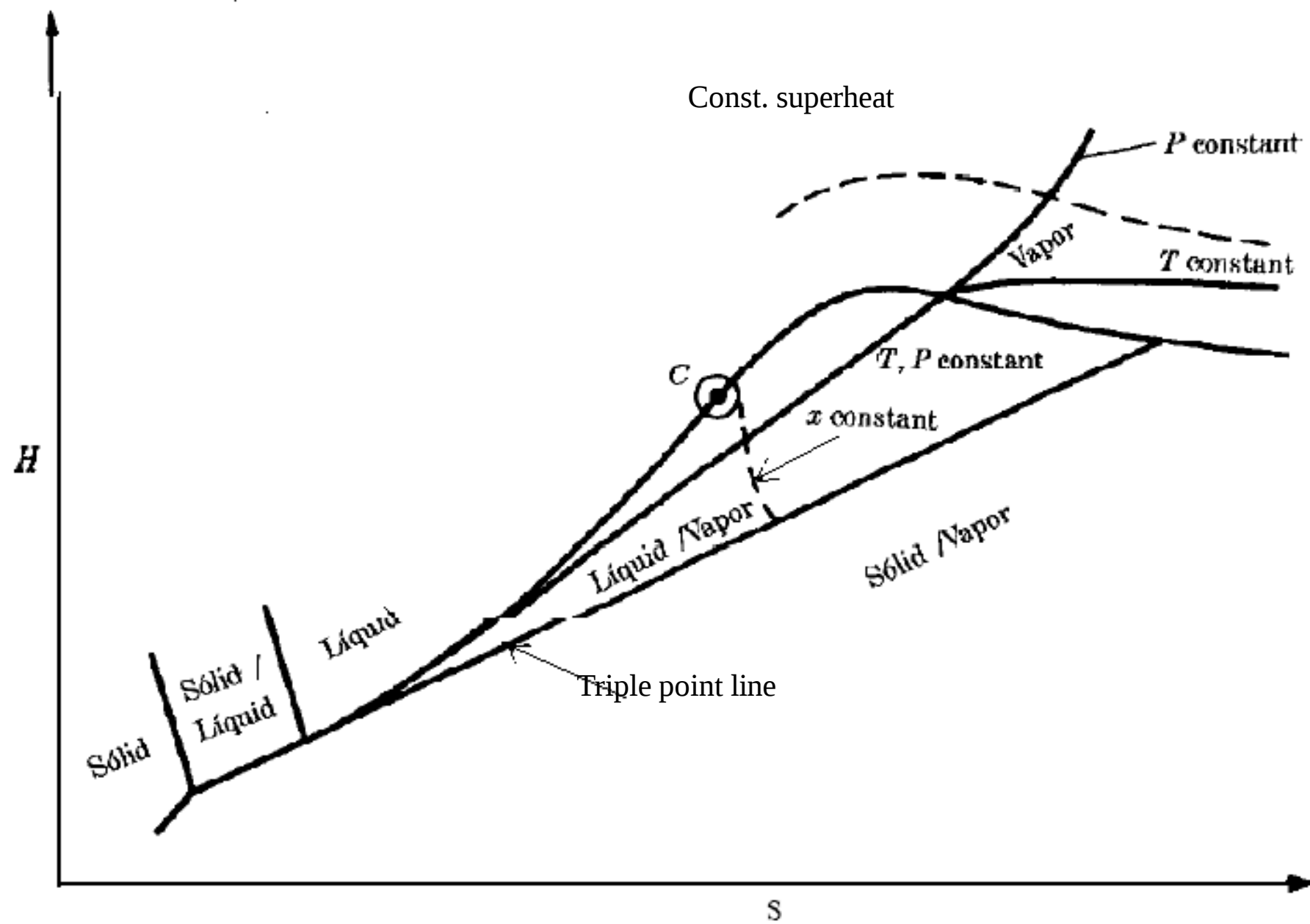
$$V = (1 - x^v)V^l + x^v V^v$$

Thermodynamic diagrams and tables

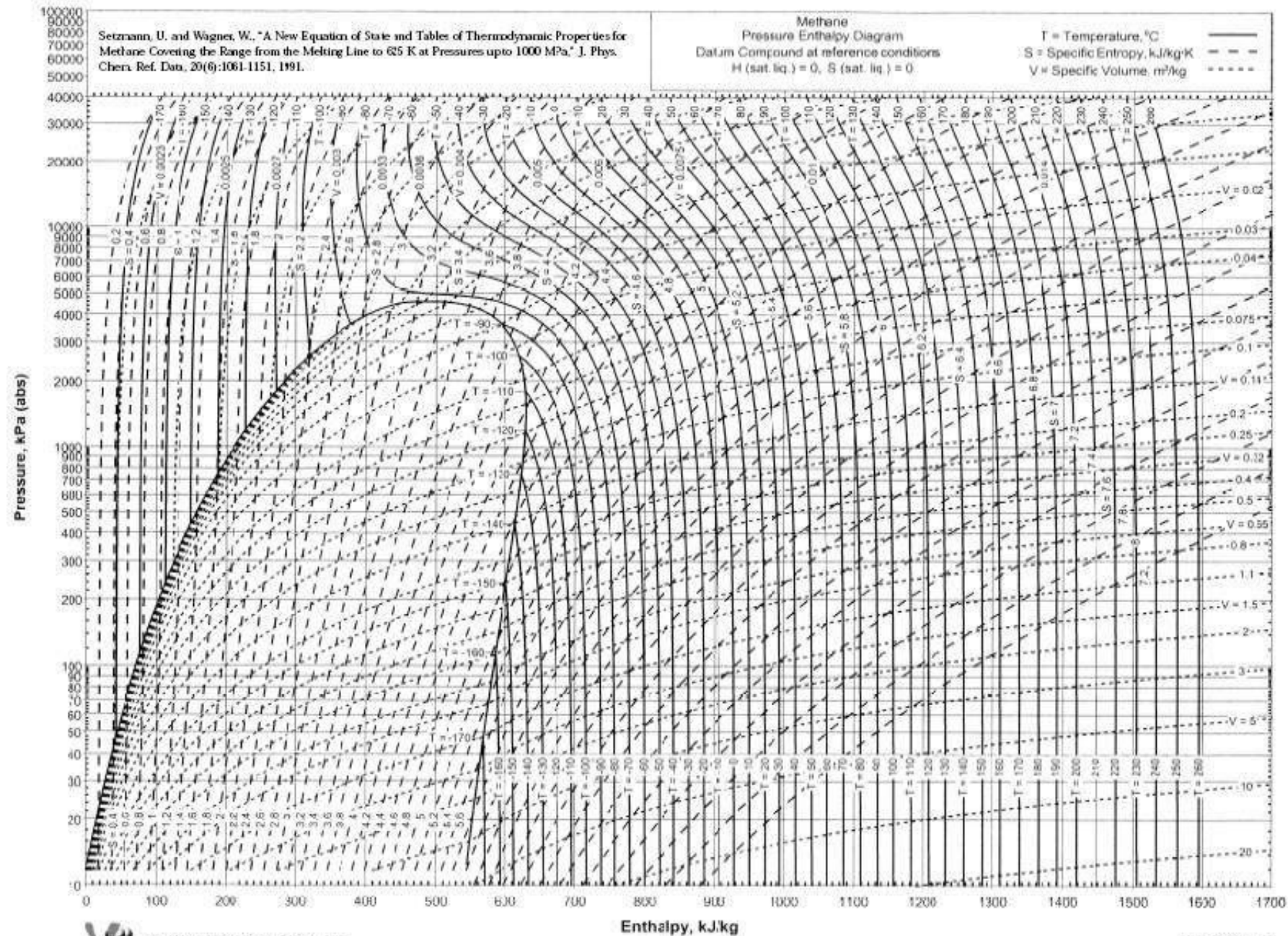
- A thermodynamic diagram represents the temperature, pressure, volume, enthalpy, and entropy of a substance on a single plot. Common diagrams are:
 - TS diagram
 - PH diagram ($\ln P$ vs. H)
 - HS diagram (Mollier diagram)
- In many instances, thermodynamic properties are reported in tables. The steam tables are the most thorough compilation of properties for a single material.

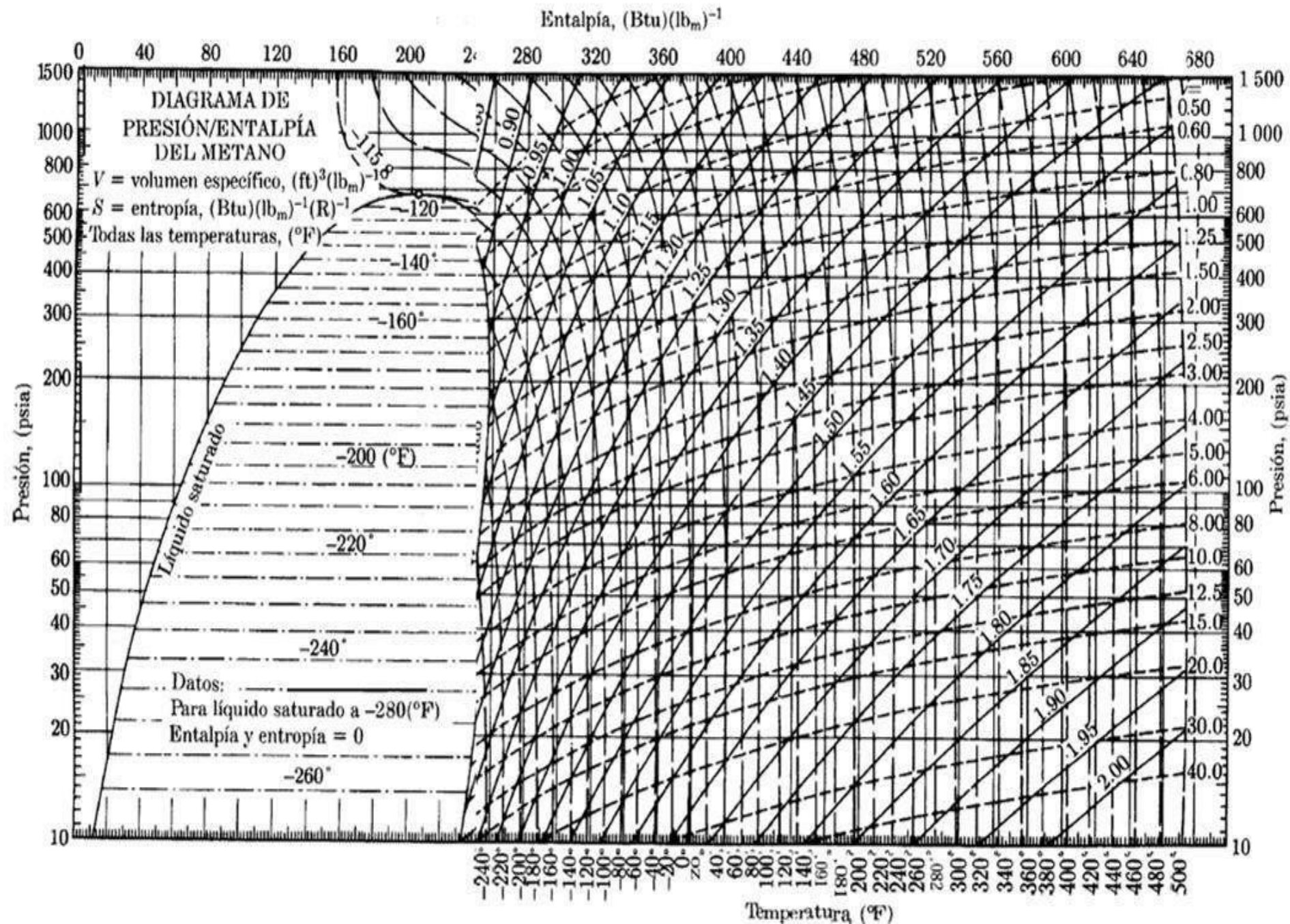


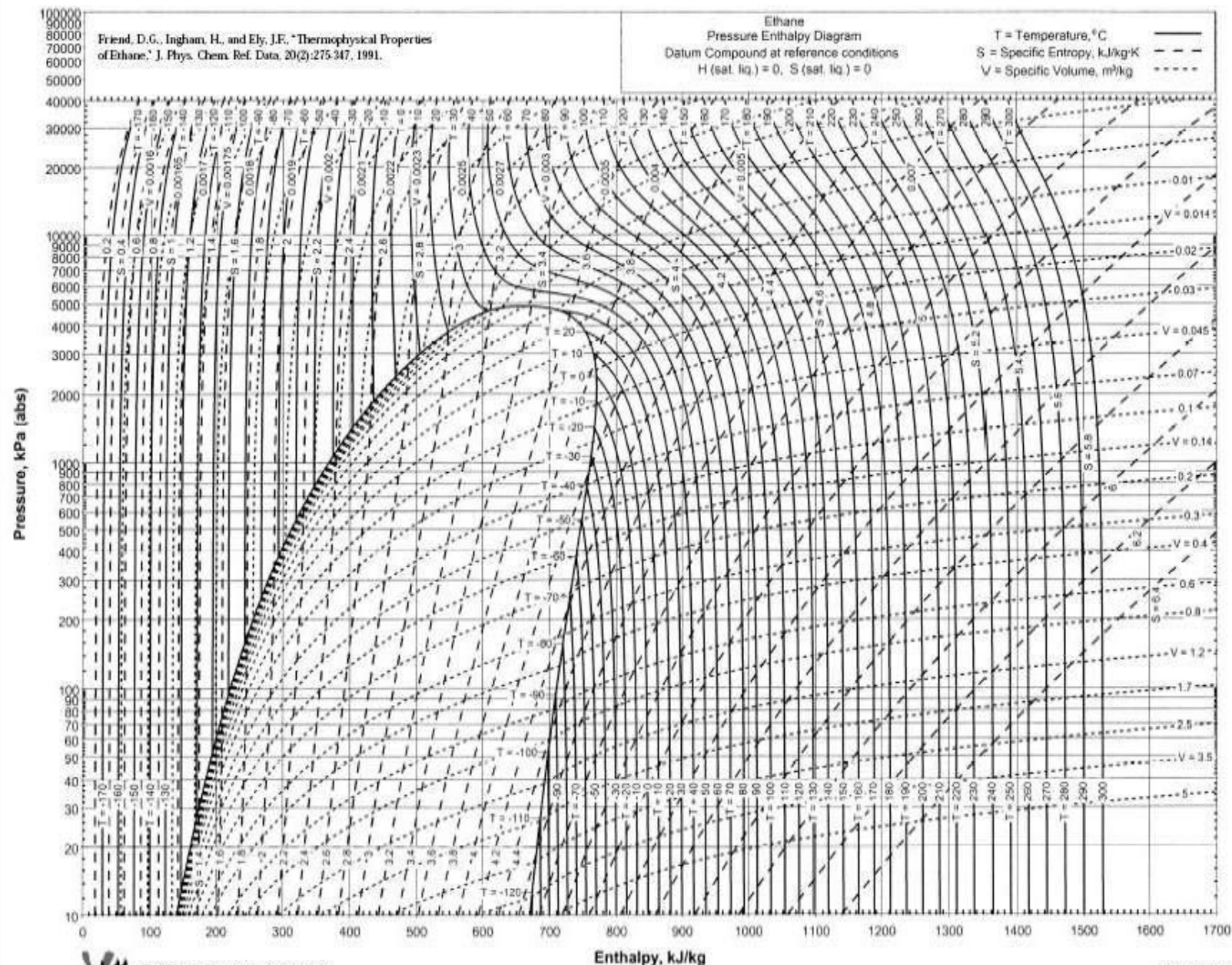




Mollier diagram



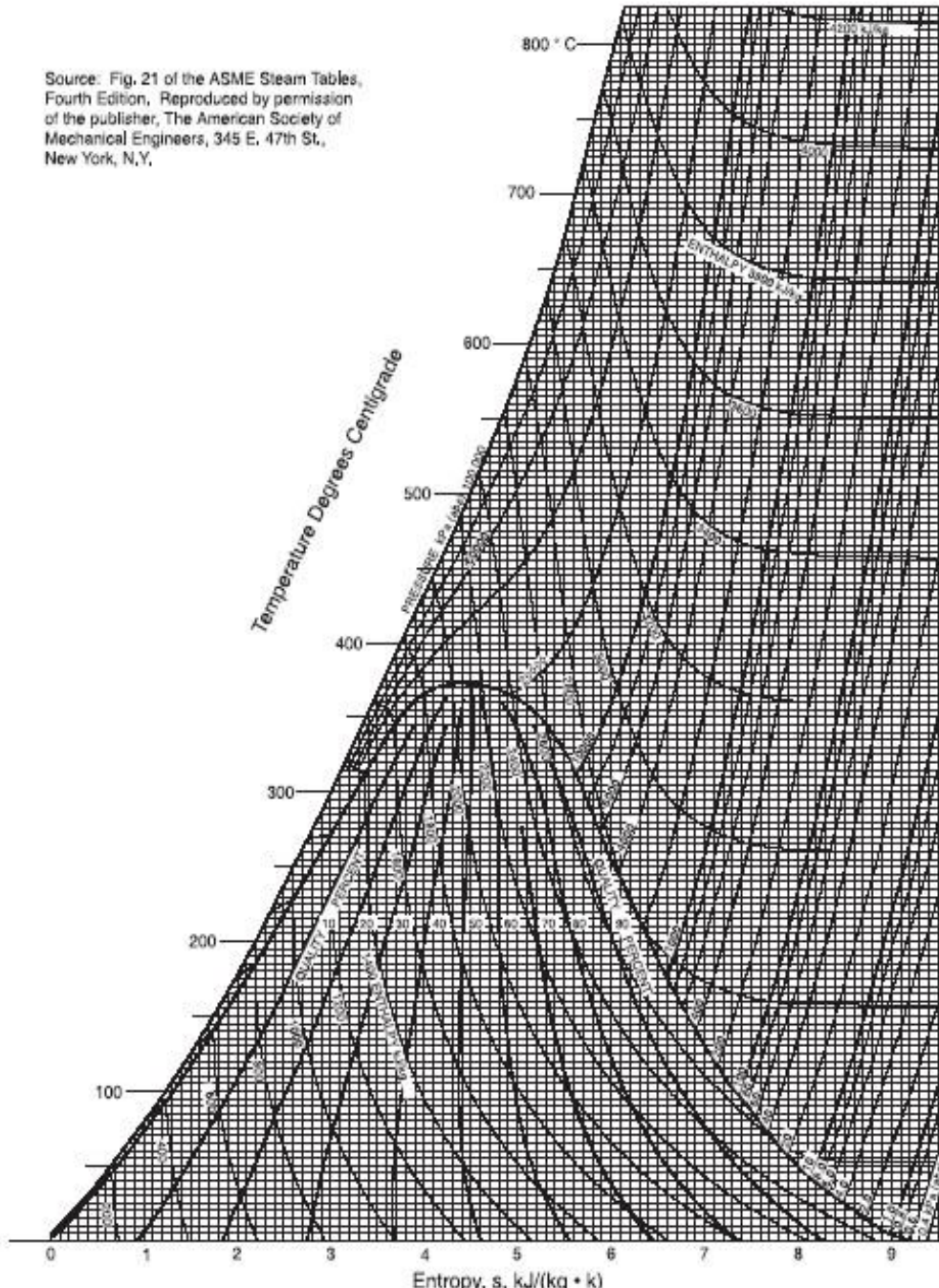




By Virtual Materials Group, Inc.

Thermodynamic Properties of Water

Source: Fig. 21 of the ASME Steam Tables, Fourth Edition, Reproduced by permission of the publisher, The American Society of Mechanical Engineers, 345 E. 47th St., New York, N.Y.



Superheated steam originally at P_1 and T_1 expands through a nozzle to an exhaust pressure P_2 . Assuming the process is reversible and adiabatic, determine the downstream state of the steam and ΔH for the following conditions:

(a) $P_1 = 1000 \text{ kPa}$, $T_1 = 250^\circ \text{C}$, and $P_2 = 200 \text{ kPa}$.

(b) $P_1 = 150 \text{ psia}$, $T_1 = 500^\circ \text{F}$, and $P_2 = 50 \text{ psia}$.

Since the process is both reversible and adiabatic, the entropy change of the steam is zero.

(a) From the steam stable and the use of interpolation,

At $P_1 = 1000 \text{ kPa}$, $T_1 = 250^\circ \text{C}$:

$$H_1 = 2942.9 \text{ kJ/kg}, S_1 = 6.9252 \text{ kJ/kg.K}$$

At $P_2 = 200 \text{ kPa}$,

$$S_2 = S_1 = 6.9252 \text{ kJ/kg.K} < S_{\text{saturated vapor}}$$

the final state is in the two-phase liquid-vapor region:

$$S_2 = (1 - x_2^v) S_2^l + x_2^v S_2^v \longrightarrow 6.9252 = 1.5301(1 - x_2^v) + 7.1268x_2^v$$

$$x_2^v = 0.9640$$

$$\Delta H = H_2 - H_1 = -315.9 \frac{\text{kJ}}{\text{kg}}$$

$$H_2 = (1 - 0.964)(504.7) + (0.964)(2706.3) = 2627.0$$

(a) From the steam stable and the use of interpolation,

At $P_1 = 150$ psia, $T_1 = 500^\circ\text{F}$:

$$H_1 = 1274.3 \text{ Btu/lbm}, S_1 = 1.6602 \text{ Btu/lbm.R}$$

At $P_2 = 50$ psia,

$$S_2 = S_1 = 1.6602 \text{ Btu/lbm.K} > S_{\text{saturated vapor}},$$

the final state is in the superheat region:

$$T_2 = 283.28^\circ\text{F} \quad H_2 = 1175.3 \text{ Btu / lbm}$$

$$\Delta H = H_2 - H_1 = -99.0 \frac{\text{Btu}}{\text{lbm}}$$

A 1.5 m³ tank contains 500 kg of liquid water in equilibrium with pure water vapor, which fills the remainder of the tank. The temperature and pressure are 100 °C, and 101.33 kPa. From a water line at a constant temperature of 70 °C and a constant pressure somewhat above 101.33 kPa, 750 kg of liquid is bled into the tank. If the temperature and pressure in the tank are not to change as a result of the process, how much energy as heat must be transferred to the tank?

Energy balance: $\frac{d(mU)_{cv}}{dt} = \dot{Q} + H'\dot{m}'$ $\rightarrow$ $\frac{d(mU)_{cv}}{dt} = \dot{Q} + H' \frac{dm_{cv}}{dt}$

$\downarrow$

$Q = \Delta(mU)_{cv} - H'\Delta m_{cv}$

$\downarrow$ $H = U + PV$

$Q = \Delta(mH)_{cv} - \Delta(PmV)_{cv} - H'\Delta m_{cv}$

$\downarrow$

$Q = (m_2 H_2)_{cv} - (m_1 H_1)_{cv} - H'\Delta m_{cv}$

At the end of the process, the tank still contains saturated liquid and saturated vapor in equilibrium at 100 °C and 101.33 kPa.

$$Q = (m_2 H_2)_{cv} - (m_1 H_1)_{cv} - H' \Delta m_{cv}$$

$$H' = 293.0 \frac{\text{kJ}}{\text{kg}} \quad \text{saturated liquid @ } 70^\circ \text{C}$$

$$H_{cv}^l = 419.1 \frac{\text{kJ}}{\text{kg}} \quad \text{saturated liquid @ } 100^\circ \text{C}$$

$$H_{cv}^v = 2676.0 \frac{\text{kJ}}{\text{kg}} \quad \text{saturated vapor @ } 100^\circ \text{C}$$

From the steam table, the specific volumes of saturated liquid and saturated vapor at 100°C are 0.001044 and $1.673 \text{ m}^3/\text{kg}$, respectively.

$$(m_1 H_1)_{cv} = m_1^l H_1^l + m_1^v H_1^v = 500(419.1) + \frac{1.5 - (500)(0.001044)}{1.673}(2676.0) = 211616 \text{ kJ}$$

$$m_2 = 500 + \frac{1.5 - (500)(0.001044)}{1.673} + 750 = m_2^l + m_2^v$$

$$1.5 = 1.673 m_2^v + 0.001044 m_2^l$$

$$(m_2 H_2)_{cv} = m_2^l H_2^l + m_2^v H_2^v = 524458 \text{ kJ}$$

$$Q = (m_2 H_2)_{cv} - (m_1 H_1)_{cv} - H' \Delta m_{cv} = 524458 - 211616 - (750)(293.0) = 93092 \text{ kJ}$$

Generalized Property Correlations for Gases (Korelasi Umum untuk Sifat- sifat Gas)

$$\frac{H^R}{RT} = -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} \quad \begin{matrix} P = P_c P_r \\ T = T_c T_r \end{matrix} \quad \frac{H^R}{RT_c} = -T_r \int_0^{P_r} \left(\frac{\partial Z}{\partial T_r} \right)_{P_r} \frac{dP_r}{P_r}$$

$$\frac{S^R}{R} = -T \int_0^P \left(\frac{\partial Z}{\partial T} \right)_P \frac{dP}{P} - \int_0^P (Z - 1) \frac{dP}{P} \quad \rightarrow \quad \frac{S^R}{R} = -T_r \int_0^{P_r} \left(\frac{\partial Z}{\partial T_r} \right)_{P_r} \frac{dP_r}{P_r} - \int_0^{P_r} (Z - 1) \frac{dP_r}{P_r}$$

$$\frac{H^R}{RT_c} = -T_r^2 \int_0^{P_r} \left(\frac{\partial Z}{\partial T_r} \right)_{P_r} \frac{dP_r}{P_r}$$

$$\frac{S^R}{R} = -T_r \int_0^{P_r} \left(\frac{\partial Z}{\partial T_r} \right)_{P_r} \frac{dP_r}{P_r} - \int_0^{P_r} (Z - 1) \frac{dP_r}{P_r}$$

$$Z = Z^0 + \omega Z^1$$

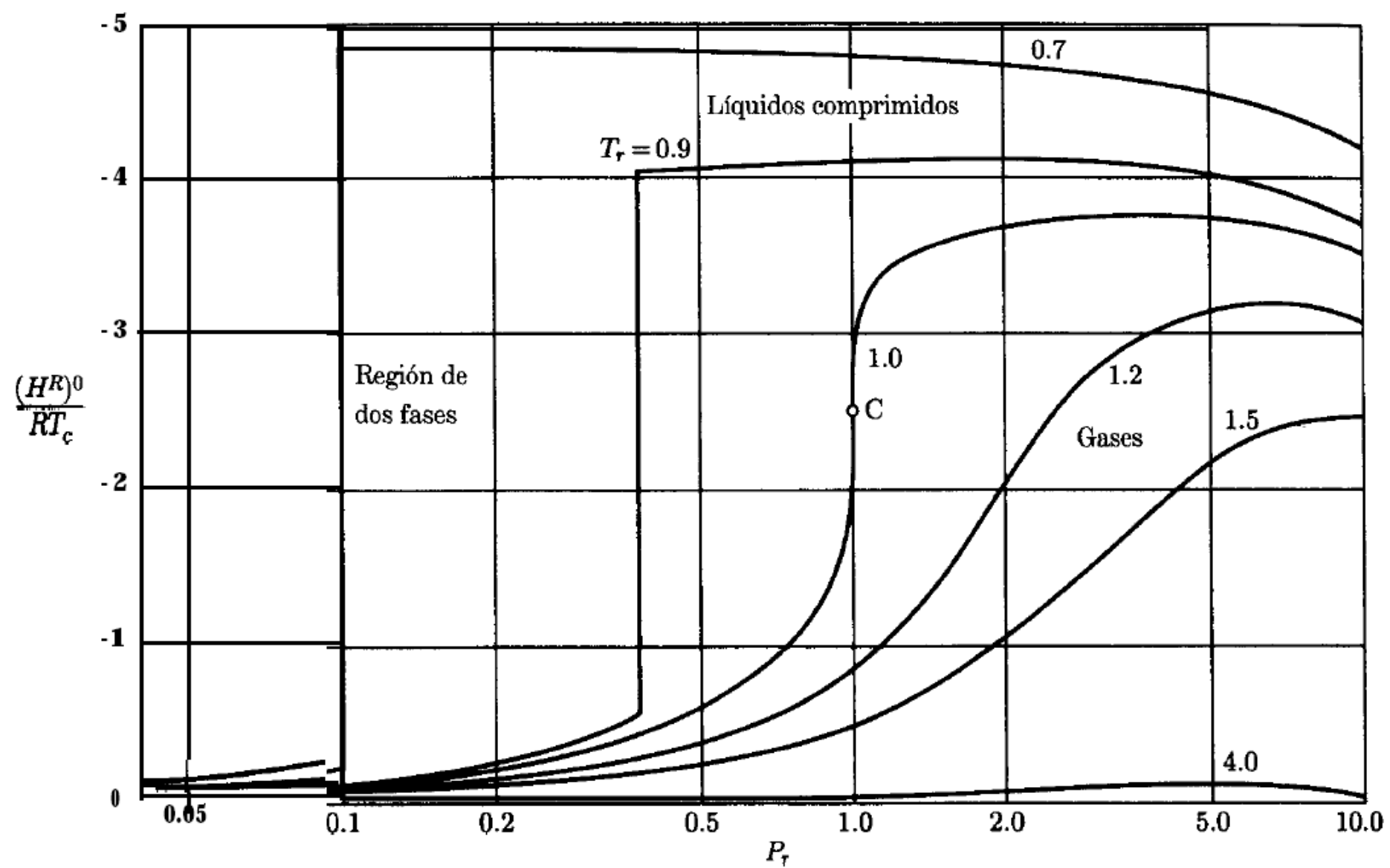
$$\frac{H^R}{RT_c} = \left[-T_r^2 \int_0^{P_r} \left(\frac{\partial Z^0}{\partial T_r} \right)_{P_r} \frac{dP_r}{P_r} \right] - \omega \left[-T_r^2 \int_0^{P_r} \left(\frac{\partial Z^1}{\partial T_r} \right)_{P_r} \frac{dP_r}{P_r} \right]$$

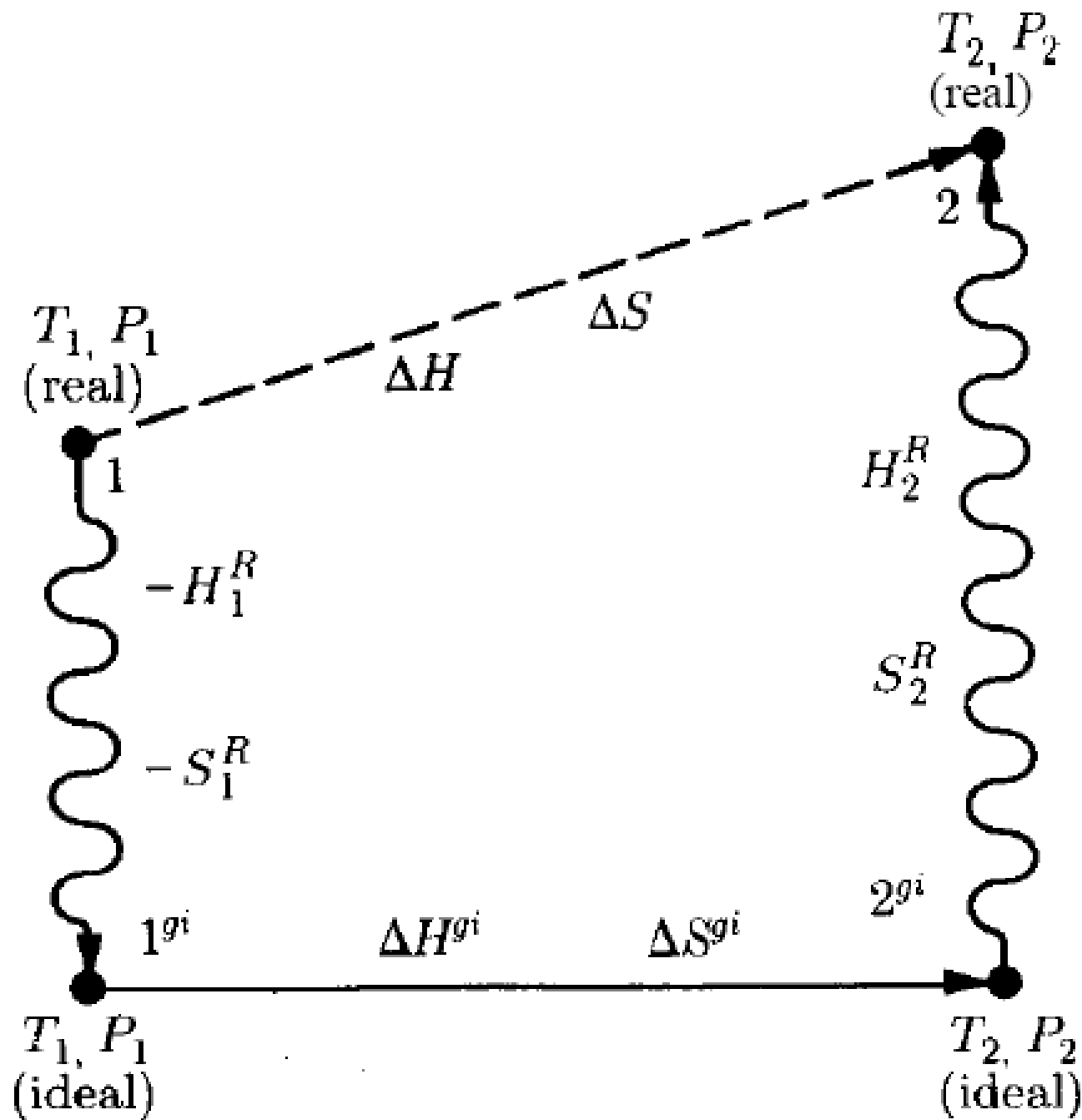
$$\frac{S^R}{R} = \left[-T_r \int_0^{P_r} \left(\frac{\partial Z^0}{\partial T_r} \right)_{P_r} \frac{dP_r}{P_r} - \int_0^{P_r} (Z^0 - 1) \frac{dP_r}{P_r} \right] - \omega \left[-T_r \int_0^{P_r} \left(\frac{\partial Z^1}{\partial T_r} \right)_{P_r} \frac{dP_r}{P_r} - \int_0^{P_r} (Z^1 - 1) \frac{dP_r}{P_r} \right]$$

$$\frac{H^R}{RT_c} = \left[\frac{(H^R)^0}{RT_c} + \omega \frac{(H^R)^1}{RT_c} \right] = HRB(TR, PR, OMEGA)$$

$$\frac{S^R}{R} = \left[\frac{(S^R)^0}{R} + \omega \frac{(S^R)^1}{R} \right] = SRB(TR, PR, OMEGA)$$

Table E5 ~ E12





$$\left. \begin{aligned} H_2 &= H_0^{ig} + \int_{T_0}^{T_2} C_P^{ig} dT + H_2^R \\ H_1 &= H_0^{ig} + \int_{T_0}^{T_1} C_P^{ig} dT + H_1^R \end{aligned} \right\} \longrightarrow \Delta H = \int_{T_1}^{T_2} C_P^{ig} dT + H_2^R - H_1^R$$

$$\downarrow$$

$$\Delta H = \langle C_P^{ig} \rangle_H (T_2 - T_1) + H_2^R - H_1^R$$

$$\left. \begin{aligned} S_2 &= S_0^{ig} + \int_{T_0}^{T_2} C_P^{ig} \frac{dT}{T} - R \ln \frac{P_2}{P_0} + S_2^R \\ S_1 &= S_0^{ig} + \int_{T_0}^{T_1} C_P^{ig} \frac{dT}{T} - R \ln \frac{P_1}{P_0} + S_1^R \end{aligned} \right\} \longrightarrow \Delta S = \int_{T_1}^{T_2} C_P^{ig} \frac{dT}{T} - R \ln \frac{P_2}{P_1} + S_2^R - S_1^R$$

$$\downarrow$$

$$\Delta S = \langle C_P^{ig} \rangle_S \ln \frac{P_2}{P_1} + S_2^R - S_1^R$$

Use of diagram for residual property calculation (Lee-Kesler Method)

- [Generalized Correlation for Compressibility factor](#)
- [Generalized Correlation for Residual Enthalpy](#)
- [Generalized Correlation for Residual Entropy](#)

Estimate V, U, H, and S for 1-butene vapor at 200 °C and 70 bar if H and S are set equal to zero for saturated liquid at 0 °C. Assume that only data available are:

$$T_c = 420K \quad P_c = 40.43 \text{ bar} \quad \omega = 0.191 \quad T_n = 266.9K \text{ (normal boiling pt.)}$$

$$C_p^{ig} / R = 1.967 + 31.630 \times 10^{-3} T - 9.837 \times 10^{-6} T^2$$

$$T_r = 1.127 \quad P_r = 1.731$$

$$\begin{aligned} Z &= Z^0 + \omega Z^1 \\ &= 0.485 + (0.191)(0.142) \\ &= 0.512 \end{aligned}$$

$$V = \frac{ZRT}{P} = 287.8 \frac{\text{cm}^3}{\text{mol}}$$

Assuming four steps:

- (a) Vaporization at T_1 and $P_1 = P^{\text{sat}}$
- (b) Transition to the ideal-gas state at (T_1, P_1)
- (c) Change to (T_2, P_2) in the ideal-gas state
- (d) Transition to the actual final state at (T_2, P_2)

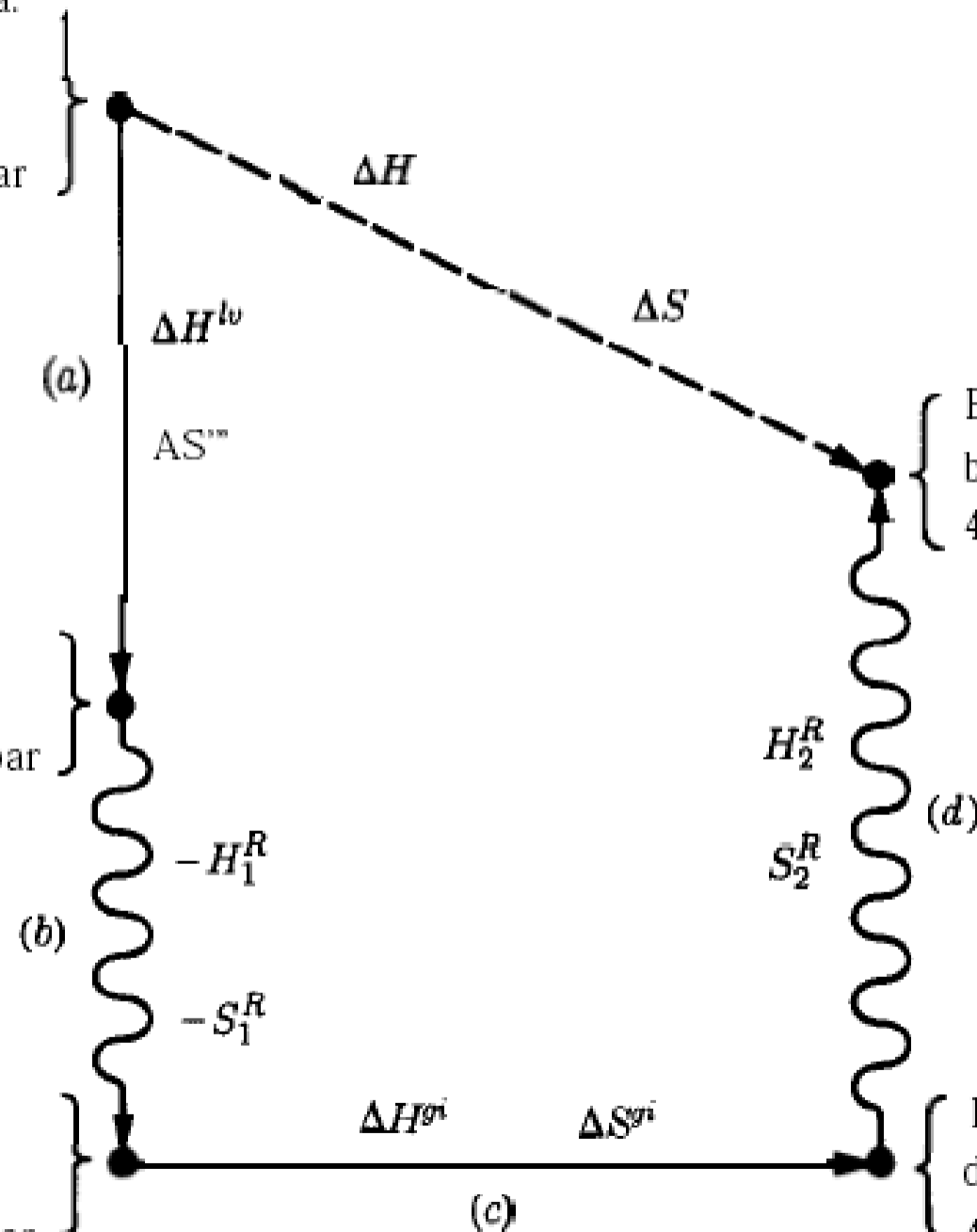
Estado de referencia:
 Buteno líquido
 saturado a
 273.15 K, 1.2771 bar

Vapor saturado
 de buteno a
 273.15 K, 1.2771 bar

Buteno en el estado
 de gas ideal a
 273.15 K, 1.2771 bar

Estado final del
 buteno a
 473.15 K, 70 bar

Buteno en el estado
 de gas ideal a
 473.15 K, 70 bar



Step (a)

$$\ln P^{sat} = A - \frac{B}{T} \quad \left\{ \begin{array}{l} \ln 1.0133 = A - \frac{B}{266.9} \\ \ln 40.43 = A - \frac{B}{420} \end{array} \right. \quad \left. \begin{array}{l} \text{For } T = 273.15\text{K}, \\ P^{sat} = 1.2771 \text{ bar} \end{array} \right.$$

The latent heat at normal boiling point:

$$\frac{\Delta H_n^{lv}}{RT_n} = \frac{1.092(\ln P_c - 1.013)}{0.930 - T_{r_n}} = 9.979$$

The latent heat at 273.15 K:

$$\frac{\Delta H^{lv}}{\Delta H_n^{lv}} = \left(\frac{1 - T_r}{1 - T_{r_n}} \right)^{0.38}$$

$$\Delta S^{lv} = 79.84 \frac{\text{J}}{\text{mol} \cdot \text{K}}$$

$$\Delta H^{\alpha\beta} = T \Delta S^{\alpha\beta}$$

$$\Delta H^{lv} = 21810 \frac{\text{J}}{\text{mol}}$$

Step (b)

$$T_r = 0.650 \quad P_r = 0.0316$$

$$\frac{H^R}{RT_c} = \left[\frac{(H^R)^0}{RT_c} + \omega \frac{(H^R)^1}{RT_c} \right] = \text{HRB}(TR, PR, \text{OMEGA}) = -0.0985$$

$$-H_1^R = 344 \frac{\text{J}}{\text{mol}}$$

$$\frac{S^R}{R} = \left[\frac{(S^R)^0}{R} + \omega \frac{(S^R)^1}{R} \right] = \text{SRB}(TR, PR, \text{OMEGA}) = -0.1063$$

$$-S_1^R = 0.88 \frac{\text{J}}{\text{mol} \cdot \text{K}}$$

Step (c)

$$\Delta H^{ig} = 8.314 \times ICPH(273.15, 473.15; 1.967, 31.630E-3, -9.837E-6, 0.0) = 20564 \frac{J}{mol}$$

$$\begin{aligned}\Delta S^{ig} &= 8.314 \times ICPS(273.15, 473.15; 1.967, 31.630E-3, -9.837E-6, 0.0) - 8.314 \ln \frac{70}{1.2771} \\ &= 22.18 \frac{J}{mol \cdot K}\end{aligned}$$

Step (d)

$$T_r = 1.127$$

$$P_r = 1.731$$

$$\frac{H^R}{RT_c} = \left[\frac{(H^R)^0}{RT_c} + \omega \frac{(H^R)^1}{RT_c} \right] = -2.430$$

$$H_2^R = -8485 \frac{J}{mol}$$

$$\frac{S^R}{R} = \left[\frac{(S^R)^0}{R} + \omega \frac{(S^R)^1}{R} \right] = -1.705$$

$$S_2^R = -14.18 \frac{J}{mol \cdot K}$$

Total

$$H = \Delta H = 21810 - (-344) + 20564 - 8485 = 34233 \frac{J}{mol}$$

$$S = \Delta S = 79.84 - (-0.88) + 22.18 - 14.18 = 88.72 \frac{J}{mol \cdot K}$$

$$U = H - PV = 34233 - (70)(287.8)/10 = 32218 \frac{J}{mol}$$

TERIMA KASIH

FOR YOUR ATTENTION

